

Recombination ("Free-Bound") ASTR 300B

The recombination rate into a level $(n, l) \leftarrow$ bound state is given by $\alpha_{ne}(T) = \langle \sigma v \rangle$ where the average $\langle \cdot \rangle$ is over $\sigma(E)$ and the distribution of speeds/energies of the electrons. For a thermal gas, this distribution is a Maxwell-Boltzmann Distribution:

$$\alpha_{ne}(T) = \int_0^\infty \sigma_{rr,ne}(E) \left(\frac{8\pi kT}{\pi m_e} \right)^{1/2} \frac{E}{kT} e^{-E/kT} \frac{dE}{kT} \text{ cm}^3 \text{ s}^{-1}$$

\circ NOTE: this is related to σ_{pi} through "Milne Relation"

See Hummer & Storey 1987 for H atom calculation of α_{ne}

There are 2 limiting cases: (Baker & Mengel 1938)

- ① CASE A - High T ($\sim 10^6$ K) shocks, collisionally ionized environments
 Optically thin to ionizing radiation \Rightarrow ionizing photons emitted during recombination escape

$$\alpha_A(T) = \sum_{n=1}^{\infty} \sum_{l=0}^{n-1} \alpha_{ne}(T)$$

- ② CASE B - HII regions, Planetary Nebulae etc.

Optically thick to radiation with $h\nu \geq 13.6 \text{ eV} \Rightarrow$ ionizing photons emitted during recombination are re-absorbed. Recombination to $n=1$ does NOT affect ionization - only recombination to $n \geq 2$ reduces/balances ionization.

$$\alpha_B(T) = \sum_{n=2}^{\infty} \sum_{l=0}^{n-1} \alpha_{ne}(T) = \alpha_A(T) - \alpha_{1s}(T)$$

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