

Blackbody Radiation

ASTR 300B

A blackbody does not reflect or scatter light but absorbs and emits light completely.

Planck's Law describes the specific intensity of a blackbody:

$$B_\nu(T) = \frac{2h\nu^3}{c^2} \cdot \frac{1}{e^{h\nu/kT} - 1} \quad \text{erg} \cdot \text{s}^{-1} \cdot \text{cm}^{-2} \cdot \text{ster}^{-1} \cdot \text{Hz}^{-1}$$

We can relate this to $B_\lambda(T)$:

$$\begin{aligned} B_\lambda(T) &= \frac{c}{\lambda^2} B_\nu(T) \\ &= \frac{c}{\lambda} \cdot \frac{2h}{c^2} \cdot \left(\frac{c^3}{\lambda^3}\right) \cdot \frac{1}{e^{hc/\lambda kT} - 1} \\ &= \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1} \quad \text{erg} \cdot \text{s}^{-1} \cdot \text{cm}^{-2} \cdot \text{ster}^{-1} \cdot \text{cm}^{-1} \end{aligned}$$

The total intensity of a blackbody is:

$$B(T) = \int_0^\infty B_\nu(T) d\nu = \frac{2h}{c^2} \int_0^\infty \frac{\nu^3 d\nu}{e^{h\nu/kT} - 1}$$

Let $x = h\nu/kT \Rightarrow \nu = \frac{kT}{h} x \Rightarrow d\nu = \frac{kT}{h} dx$

$$B(T) = \frac{2h}{c^2} \cdot \left(\frac{kT}{h}\right)^4 \int_0^\infty \frac{x^3 dx}{e^x - 1}$$

← This integral is $\int_0^\infty \frac{x^a dx}{e^x - 1} = \Gamma(a+1) \zeta(a+1)$

REIMANN Zeta Fctn. ↓
Gamma Fctn. ↓

$$B(T) = \frac{2}{c^2} \cdot \frac{k^4 T^4}{h^3} \cdot \Gamma(4) \cdot \zeta(4) = \frac{2}{c^2} \frac{k^4 T^4}{h^3} \cdot 6 \cdot \frac{\pi^4}{90}$$

$$= \frac{2\pi^4 k^4}{15 c^2 h^3} T^4$$

← NOTICE T^4 dependence
 $\text{erg} \cdot \text{s}^{-1} \cdot \text{cm}^{-2} \cdot \text{ster}^{-1}$

Emergent Flux (or "astrophysical flux") of a black body is :

$$F^+ = \pi \cdot B(T) = \pi \cdot \frac{2\pi^4 k^4}{15 c^2 h^3} T^4$$

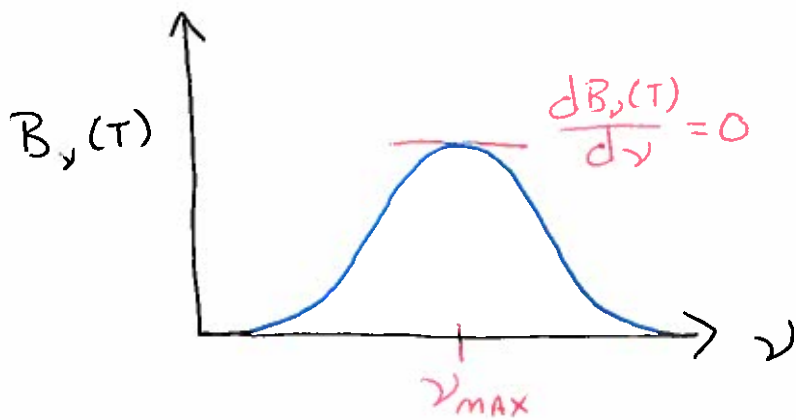
$$= \sigma T^4 \quad \text{erg s}^{-1} \cdot \text{cm}^{-2}$$

σ = Stephan-Boltzmann Constant $\sim 5.67 \times 10^{-5} \text{ erg s}^{-1} \cdot \text{cm}^{-2} \cdot \text{K}^{-4}$

For a star (\approx a good blackbody!) :

$$L_* = 4\pi R_*^2 F_*^+ = 4\pi R_*^2 \sigma T_*^4$$

Let plot the Planck Function :



The Planck function goes to zero at the limits $\rightarrow 0$ and $\rightarrow \infty$ and has a single maximum. We have to use numerical techniques to find the maximum :

$$\frac{dB_\nu(T)}{d\nu} = 0 \quad \Rightarrow \quad \frac{\nu_{\text{MAX}}}{T} = 5.879 \times 10^{10} \text{ Hz} \cdot \text{K}^{-1}$$

$$\frac{dB_\lambda(T)}{d\nu} = 0 \quad \Rightarrow \quad \lambda_{\text{MAX}} T = 0.2898 \text{ cm} \cdot \text{K}$$

Wein's Displacement Law

LIMITS of the Planck Function

Notice that $h\nu/k_B T$ has units of T (kelvin)

① High ν limit: $\frac{h\nu}{k_B T} \gg 1$ Wein Approximation

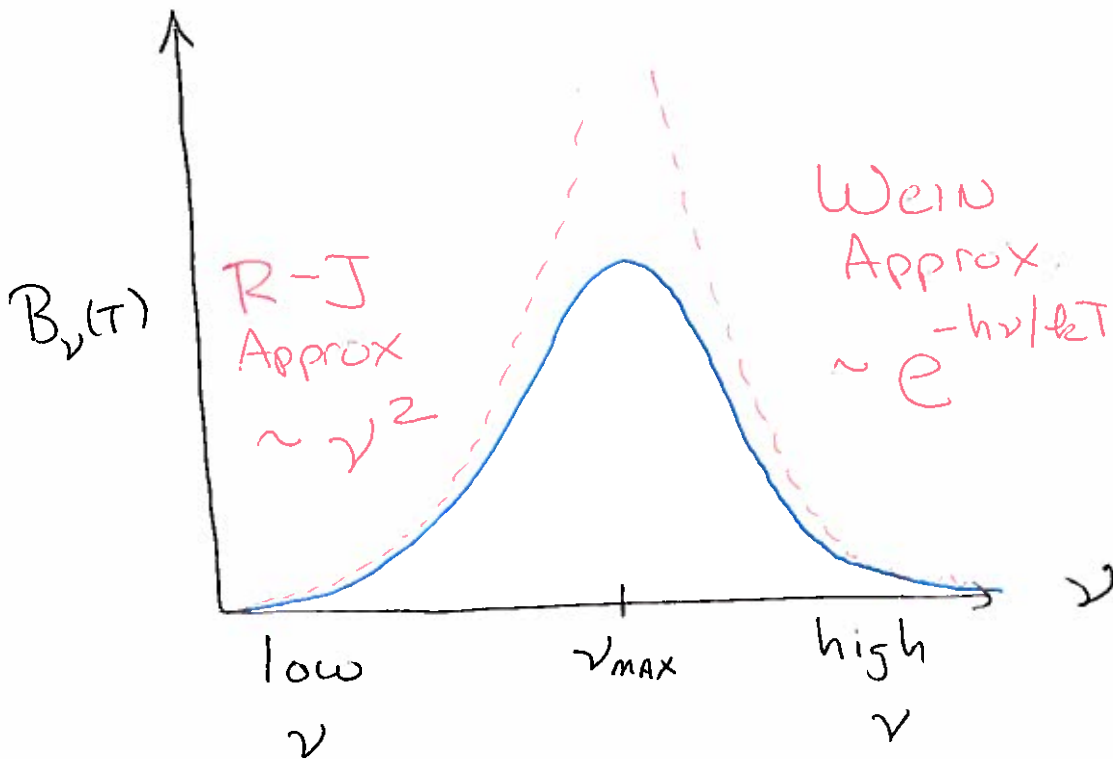
$\lim_{\nu \rightarrow \text{large}} B_\nu(T) \rightarrow e^{-h\nu/k_B T}$ exponential dependence at high ν

② Low ν limit: $\frac{h\nu}{k_B T} \ll 1$ Rayleigh-Jeans Approximation

Taylor expand: $e^{h\nu/k_B T} \approx 1 + \frac{h\nu}{k_B T}$

$\lim_{\nu \rightarrow \text{small}} B_\nu(T) = \frac{2h\nu^3}{c^2} \cdot \frac{1}{1 + \frac{h\nu}{k_B T} - 1} = \frac{2k_B T}{c^2} \nu^2$

Intensity linearly proportional to T
Notice ν^2 dependence at low ν



Energy Density of a Blackbody

$$u_\nu = \frac{1}{c} \oint B_\nu(T) d\Omega$$

Assuming $B_\nu(\theta, \phi) = \text{const}$

$$u_\nu = \frac{4\pi}{c} \cdot \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT} - 1} = \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{h\nu/kT} - 1}$$

$\text{erg} \cdot \text{cm}^{-3} \cdot \text{Hz}^{-1}$

The total energy density is given by:

$$\begin{aligned} u &= \frac{4\pi}{c} \cdot \int_0^\infty B_\nu(T) d\nu = \frac{4\pi}{c} \cdot \frac{\sigma T^4}{\pi} \\ &= \frac{4\sigma}{c} T^4 \\ &= a T^4 \quad \text{erg} \cdot \text{cm}^{-3} \end{aligned}$$

called the radiation constant

$$a = 7.567 \times 10^{-15} \text{ erg} \cdot \text{cm}^{-3} \cdot \text{K}^{-4}$$