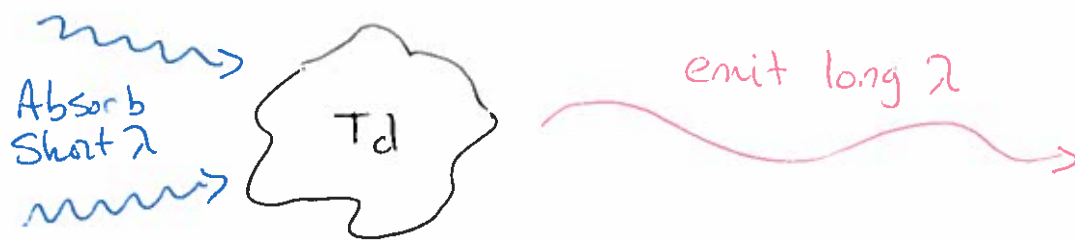


DUST EMISSION

ASTR 300B



CONSIDER PURE EMISSION (AND IGNORE ANY BACKGROUND):

$$I_\nu = \cancel{I_\nu(0)} e^{-\tau_\nu} + S_\nu (1 - e^{-\tau_\nu})$$

↑ IGNORE

Thermal emission from dust $\Rightarrow S_\nu = B_\nu(T_d) = \frac{j_\nu}{\alpha_\nu}$

$$j_\nu = B_\nu(T_d) \alpha_\nu$$

If emission is in optically thin limit:

$$I_\nu = B_\nu(T_d) \tau_\nu$$

$$\tau_\nu = \int_0^L \alpha_\nu ds = \int_0^L n_d \sigma_\nu^{\text{dust}} ds = \int_0^L \rho_d \chi_\nu^{\text{dust}} ds = \int_0^L \rho_{\text{gas}} \chi_\nu^{\text{gas}} ds$$

$\frac{\text{cm}^{-3} \cdot \text{cm}^2}{\text{cm}^3 \cdot \frac{\text{cm}^2}{\text{g of dust}}}$

Remember $\frac{\rho_{\text{gas}}}{\rho_{\text{dust}}} = \left\langle \frac{\text{mass in gas}}{\text{mass in dust}} \right\rangle \sim 100 \Rightarrow$ We can relate τ_ν to how much GAS along line of sight!

$$\chi_\nu^{\text{dust}} \sim \chi_\nu^{\text{gas}} \cdot 100$$

$$\frac{\text{cm}^2}{\text{g of dust}} \sim \frac{\text{cm}^2}{\text{g of gas}} \times \frac{100 \text{ g of gas}}{1 \text{ g of dust}}$$

PUTTING THIS ALL TOGETHER

$$\begin{aligned}
 I_\nu &= B_\nu(T_d) \cdot \int_0^L \rho_{\text{gas}} \cdot \kappa_\nu^{\text{per gas}} ds \\
 &= B_\nu(T_d) \cdot \int_0^L \mu m_H n_{\text{gas}} \cdot \frac{\kappa_\nu^{\text{dust}}}{100} ds \\
 &= B_\nu(T_d) \mu m_H \frac{\kappa_\nu^{\text{dust}}}{100} \int_0^L n_{\text{gas}} ds
 \end{aligned}$$

If we measure n_{gas} in terms of H_2 molecules then

$$I_\nu = B_\nu(T_d) \mu_{\text{H}_2} m_H \frac{\kappa_\nu^{\text{dust}}}{100} N_{\text{H}_2} \leftarrow \begin{array}{l} \text{Column Density} \\ \text{of H}_2 \end{array}$$

cm^{-2}

$\frac{I_\nu^{\text{dust emission}}}{\nu} \sim N_{\text{H}_2}$
optically thin limit

NOTE: We have assumed $T_d = \text{const}$ along line of sight
and that $\kappa_\nu^{\text{dust}} = \text{const}$ along line of sight

If T_d, κ_ν vary - have to go back and solve radiative transfer equation:

$$\frac{dI_\nu}{ds} = j_\nu = \alpha_\nu B_\nu(T_d(s)) = \frac{\mu m_H}{100} n(s) \kappa_\nu^{\text{dust}}(s) B_\nu(T_d(s))$$

AND INTEGRATE ...
$$I_\nu = \frac{\mu m_H}{100} \int_0^L n(s) \kappa_\nu^{\text{dust}}(s) B_\nu(T_d(s)) ds$$

We can relate ρ_{gas} to n_{gas} through the mean molecular weight:

There are different definitions depending on whether you are calculating per H_2 molecule or per "particle":

$$\mu_{\text{H}_2} m_{\text{H}} N(\text{H}_2) = M$$

Number of H_2 molecules
($\times 2$ column density)

Typical
Abundance Ratios
in ISMs

$$M = M(\text{H}) + M(\text{He}) + M(\text{Z})$$

$$\frac{M(\text{H})}{M} \sim 0.71 \quad \frac{M(\text{He})}{M} \sim 0.27$$

$$\frac{M(\text{Z})}{M} \sim 0.02 \quad \frac{N(\text{H})}{N(\text{He})} \sim 10$$

$$\mu_{\text{H}_2} = \frac{M}{m_{\text{H}} N(\text{H}_2)} = \frac{2M}{m_{\text{H}} N(\text{H})} = \frac{2M}{M(\text{H})} \approx 2.8$$

Per free particle:

$$\mu_{\text{p}} \approx \frac{M}{m_{\text{H}} [N(\text{H}_2) + N(\text{He}) + N(\text{Z})]}$$

$$= \frac{M}{m_{\text{H}} \left[\frac{N(\text{H})}{2} + \frac{N(\text{H})}{10} + \text{negligible} \right]}$$

$$= \frac{M}{\frac{M(\text{H})}{2} + \frac{M(\text{H})}{10}} = \frac{M}{M(\text{H})} \cdot \frac{5}{3} \approx 2.34$$