

Boltzmann Equation

ASTR 300B

The Boltzmann Equation give the distribution of states for a system with total energy E at a temperature T .

$$\frac{n_2}{n_1} = \frac{g_2}{g_1} e^{-\Delta E_{21}/kT_{ex}}$$

$\Delta E_{21} = E_2 - E_1$
 $k \leftarrow$ Boltzmann's Constant
 $k \approx 1.381 \times 10^{-16} \text{ erg K}^{-1}$
 $T_{ex} \leftarrow$ excitation temperature

Statistical (degeneracies of weights (energy levels)) For instance, e^- with angular momentum l have $g = (2l+1)$

We shall now derive this expression.

ASSUME (1) Total Energy CONSTANT

$$E = \sum_{i=1}^{\infty} n_i \epsilon_i \quad n_i = \text{number of particles with energy } \epsilon_i$$

(2) Total number of particles is CONSTANT (AND LARGE)

$$N = \sum_i n_i$$

(3) Particles are distinguishable (NOTE this may not be the case for some quantum gases)

We are going to analyze the "macrostate" Ω of a "microcanonical ensemble" of particles with energies ϵ_i

What the heck does that mean?? Let's use an example to derive the possible combinations:

Assume you have 3 CONTAINER (A state that the particle can be put into) that hold the first 9 letters of the alphabet:

a b c d e f g h i
| _ _ | | _ _ _ | | _ _ _ |

For the 1st slot, there are 9 possibilities

For the 2nd slot, there are 8 possibilities

For the 3rd slot, there are 7 possibilities

etc.

⇒ there are $9!$ possible arrangements if the containers are ignored.

BUT, containers matter! We don't care about the ordering of letters within the container, so we divide out by the arrangement possibilities within each container

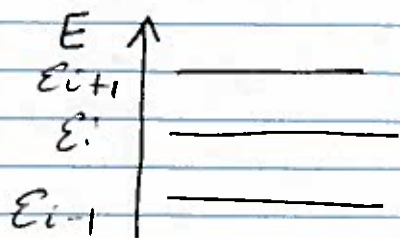
$$\Omega = \frac{9!}{2!3!4!} = 1260 \text{ ways of partitioning 9 letters into 3 containers}$$

IN General, the number of distinct arrangements of N particles into n groups containing n_1, n_2, \dots, n_n objects is

$$\Omega = \frac{N!}{n_1! n_2! \dots n_n!} = \frac{N!}{\prod_i n_i!}$$

There is one additional complication.

Each container is a state that each particle can be put into



g_i degenerate states into which n_i particles can be arranged.
 n independent energy levels

Think of this as a system of n containers with g_i subcontainers

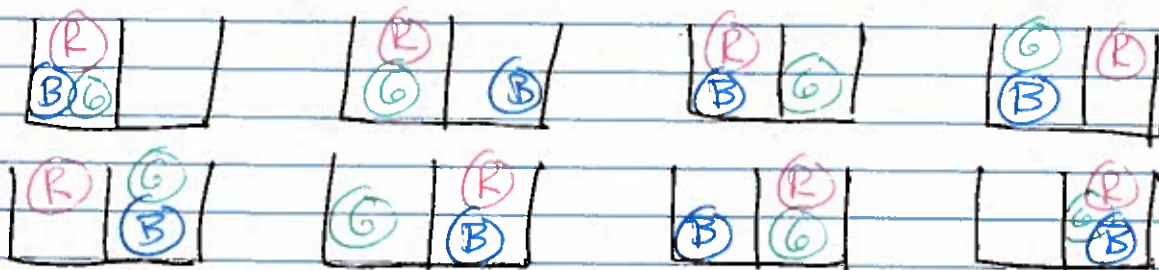
Classical case: particles in subcontainers are distinguishable.

EXAMPLE: CONSIDER 2 subcontainers w/ 3 particles



(R) (G) (B)

How many arrangements of n_i particles in g_i subcontainers are possible?



8 arrangements or $2^3 \Rightarrow g_i^{n_i}$

The total number of arrangements is then

$$\Omega = \frac{N!}{\prod_i n_i!} \prod_i g_i^{n_i}$$

The most likely macrostate MAXIMIZES Ω
 while holding N, E CONSTANT.
 $\Rightarrow d\Omega = 0$

$$d(\ln \Omega) = \frac{1}{\Omega} d\Omega = 0 \Rightarrow \text{we can also maximize } \ln \Omega$$

$$\ln \Omega = \ln N! + \sum_i n_i \ln g_i - \sum_i \ln n_i!$$

For large N , use Stirling's Formula (Approximation)

$$\ln N! = N \ln N - N$$

$$\begin{aligned} \ln \Omega &\approx N \ln N - N + \sum_i n_i \ln g_i - \sum_i n_i \ln n_i + \sum_i n_i \\ &\approx N \ln N + \sum_i n_i \ln \left(\frac{g_i}{n_i} \right) \end{aligned}$$

$$\begin{aligned} d(\ln \Omega) &= \overset{N=\text{CONST}}{0} + \sum_i dn_i \ln \left(\frac{g_i}{n_i} \right) + n_i \cdot \frac{dn_i}{g_i} - \frac{g_i}{n_i^2} dn_i \\ &= \sum_i dn_i \ln \left(\frac{g_i}{n_i} \right) - \sum_i dn_i \quad \leftarrow \text{because } N=\text{CONST} \\ &= 0 \quad \text{to maximize } d(\ln \Omega) \end{aligned}$$

Now energy conservation $\Rightarrow \sum_i \epsilon_i dn_i = 0$

So we use the method of Lagrange Multipliers to add the constraints of $dN = dE = 0$

$$\alpha \sum_i dn_i = 0$$

$$- \beta \sum_i \epsilon_i dn_i = 0$$

$$\Rightarrow \sum_i \left[\ln \left(\frac{g_i}{n_i} \right) + \alpha - \beta \epsilon_i \right] dn_i = 0$$

The only way this is true is if

$$\ln\left(\frac{g_i}{n_i}\right) + \alpha - \beta \epsilon_i = 0$$

$$\Rightarrow n_i = g_i e^{\alpha} e^{-\beta \epsilon_i}$$

$$\text{Now } N = \sum_i n_i = \sum_i g_i e^{\alpha} e^{-\beta \epsilon_i}$$

But e^{α} is a CONST, so can take out of \sum_i

$$\Rightarrow e^{\alpha} = \frac{N}{\sum_i g_i e^{-\beta \epsilon_i}} = \frac{N}{Q} \leftarrow \text{partition function}$$

$$\Rightarrow n_i = g_i \frac{N}{Q} e^{-\beta \epsilon_i}$$

* Note β has to have units of erg⁻¹

Boltzmann proved that $\beta = \frac{1}{kT}$ where

$k \approx 1.381 \times 10^{-16} \text{ erg} \cdot \text{K}^{-1}$ is Boltzmann's CONSTANT.

$$\frac{n_i}{N} = \frac{g_i}{Q} e^{-\epsilon_i/kT} \leftarrow$$

IN General
 $T = T_{\text{ex}} = \text{excitation temperature}$

$$\text{OR } \frac{n_i}{n_j} = \frac{g_i}{g_j} e^{-(\epsilon_i - \epsilon_j)/kT}$$

When all levels are "thermalized"
 $T_{\text{ex}} = T_k$
gas kinetic temperature of the particles