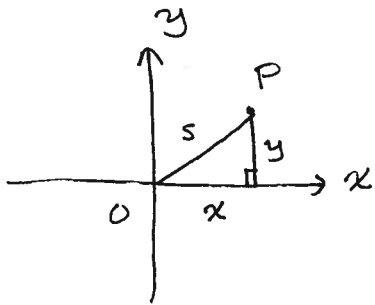


Metrics

1.

A mathematical description for how to measure distances



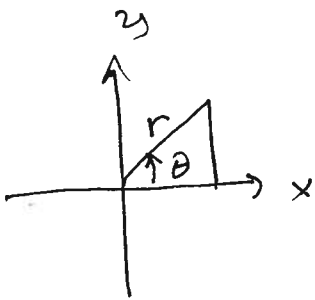
The distance between 0 and P
is $s^2 = x^2 + y^2$ Pythagorean Theorem.

In the differential limit

$$ds^2 = dx^2 + dy^2$$

coefficients are 1 \Rightarrow space is flat or Euclidean.

\mathbb{R}^2 \leftarrow 2 dimension Real space



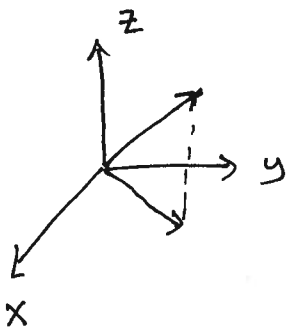
we can transform to polar coordinates

$$x = r \cos \theta \quad y = r \sin \theta$$

$$ds^2 = dr^2 + r^2 d\theta^2$$

coefficients are not 1 but this is still a flat space because a transformation exists to a metric where the coefficients are 1.

Let's extend this concept to \mathbb{R}^3 :



$$ds^2 = dx^2 + dy^2 + dz^2$$

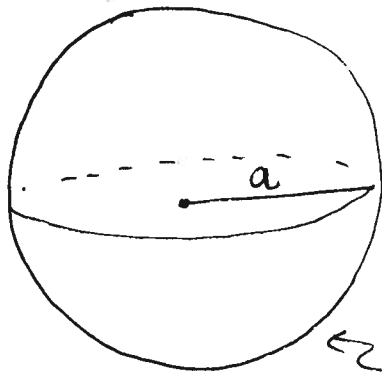
OR in spherical coordinates

$$ds^2 = dr^2 + r^2(d\theta^2 + \cos^2 \theta d\phi^2)$$

Now consider a \mathbb{R}^2 space curved in \mathbb{R}^3 :

(2.)

e.g. a Sphere



$$ds^2 = a^2 (d\theta^2 + \cos^2 \theta d\phi^2)$$



This is a scale factor

← Fixed sphere w/ radius $a \Rightarrow dr=0, r=a.$

We need to be careful when talking about curvature in \mathbb{R}^3 .

There are two types of curvature:

extrinsic = curvature of a lower dimensional space embedded in a higher dimension

intrinsic = curvature that is a fundamental property of the space (i.e. no higher dimensional embedding required)

~~The scale factor~~

Example of extrinsic curvature:

- ① take a flat sheet of paper (\mathbb{R}^2 flat)
- ② roll it into a cylinder



This \mathbb{R}^2 space has extrinsic curvature in \mathbb{R}^3 but does not have intrinsic curvature.

The spherical surface in \mathbb{R}^3 has intrinsic curvature.

What about space-time?

A flat space-time metric (Special Relativity)

$$ds^2 = \underbrace{dx^2 + dy^2 + dz^2}_{\text{Space-part}} - \underbrace{c^2 dt^2}_{\text{time-part}}$$

time distinct from spatial coordinates.

How does all of this fit together?

Lemaître-Friedmann, Robertson-Walker Metric (FRW) (General Relativity)

For an isotropic and homogeneous universe (Cosmological Principle)

$$ds^2 = -c^2 dt^2 + a^2(t) \left[\frac{dr^2}{1 - kr^2} + r^2 (d\theta^2 + \cos^2 \theta d\phi^2) \right]$$

↑ co-moving coordinates
 ↑ curvature (intrinsic)
 $k = \begin{cases} > 0 & \text{positive (e.g. "spherical")} \\ 0 & \text{Flat} \\ < 0 & \text{negative (e.g. "hyperbolic")} \end{cases}$
 ↑ scale factor of expansion (or contraction)

$$\frac{d(t)}{d(t_0)} = \frac{\text{co-moving distance measured @ time } t}{\text{co-moving distance measured today}} = \frac{a(t)}{a(t_0)} = \frac{\lambda_{em}}{\lambda_{obs}}$$

Where have we seen this before?

$$1 + z = \frac{\lambda_{obs}}{\lambda_{em}} = \frac{a(t_0)}{a(t)} = \frac{a_0}{a(t)}$$

Don't forget Hubble Law:

$$v = H \cdot d$$

$$v = \dot{d}(t) = \frac{d}{dt} \left(\frac{a(t)}{a_0} \right) d(t_0)$$

$$= \frac{\dot{a}(t)}{a_0} d(t_0)$$

$$\text{so } H(t) = \frac{\dot{d}(t)}{d(t)} = \frac{\dot{a}(t)}{a(t)} \sim \text{rate of change of scale factor.}$$

How does $a(t)$ vary w/ cosmological evolution?

General Relativity \rightarrow Einstein Field Equations $\xrightarrow{\text{solutions}}$ differential equations $\rightarrow a(t)$

$$\text{e.g. } \frac{\ddot{a}(t)}{a(t)} = -\frac{4\pi G}{3} \left(\rho + \frac{3p}{c^2} \right) + \frac{\Lambda c^2}{3}$$

Solutions of this equation are beyond what we want to do now.

$$\text{example: } E \sim h\nu \sim \frac{hc}{\lambda} \Rightarrow E \propto a^{-1}$$

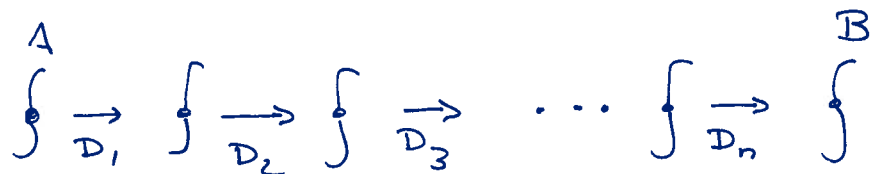
$$n_\gamma = \frac{N_\gamma}{V} \propto a^{-3} \text{ if } N_\gamma \text{ conserved.}$$

Cosmological Distance Definitions

Remember :
$$\frac{\lambda_{\text{obs}}}{\lambda_{\text{em}}} = 1+z = \frac{R(t_0)}{R(t)}$$

For $z \ll 1$ we can use $v = c \cdot z$ $D = \frac{v}{H(t)} = \frac{cz}{H(t)}$

Imagine two galaxies A and B separated by a large distance ($z \ll 1$). In order to measure the co-moving radial distance, we need a chain of observers



each observer in a galaxy whose distance is small such that $v \sim c \cdot z$ is appropriate. They must all measure D_1, D_2, \dots, D_n at the same proper time to determine

$$D_{\text{now}} = D_1 + D_2 + \dots + D_n$$

This is impractical !
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Therefore, alternative distance definitions have been devised...

(2)

(1) Co-moving radial distance

D_{now}

← what we have been using to date

(2) Angular diameter distance

$$D_A \equiv \frac{\text{size of object}}{\text{angular extent of object in the sky}} = \frac{R}{\theta}$$

(Problem: Need to know size of object!)

(3) Luminosity Distance

$$F = \frac{L}{4\pi D_L^2}$$

(Problem: Need to know L) * useful for "standard" candles!

(4) Light travel time distance

$$D_{\text{ltt}} = c \cdot (t_0 - t_{\text{em}})$$

↑
age of universe

Event Horizon is

$$D_{\text{eh}} = ct_0$$

(Need to know t_0 and figure out t_{em})

(3)

It turns out D_A & D_L are related:

$$L = 4\pi R^2 \sigma T_{em}^4 \quad \text{for a blackbody at redshift } z$$

$$T_{obs} = \frac{T_{em}}{1+z} \quad \leftarrow \text{you will prove this in Hmwk \# 8}$$

$$F_{obs} = \theta^2 \cdot \sigma T_{obs}^4$$

$$D_L^2 = \frac{L}{4\pi F_{obs}} = \frac{4\pi R^2 \sigma T_{em}^4}{4\pi \theta^2 \sigma T_{obs}^4}$$

remember $D_A = \frac{R}{\theta}$

$$D_L^2 = D_A^2 \cdot \frac{T_{em}^4}{T_{obs}^4} = D_A^2 \cdot (1+z)^4$$

$$D_L = D_A \cdot (1+z)^2$$