

Ro-vibrational Transitions

ASTR
300B

When transitions happen between 2 vibrational states of a molecule, the rotational state can also change - these are "ro-vibrational" transitions.

For a harmonic oscillator potential the solutions to Schrodinger's Equations are

$$\psi \sim H_n(r) e^{-r^2/2}$$

Hermite polynomials

$$E_v = \hbar \omega_0 (v + \frac{1}{2})$$

Vibrational quantum number

$$\omega_0 = \sqrt{\frac{k}{\mu}} \leftarrow \text{spring constant} \quad = 2\pi\nu_0 \leftarrow \text{reduced mass}$$

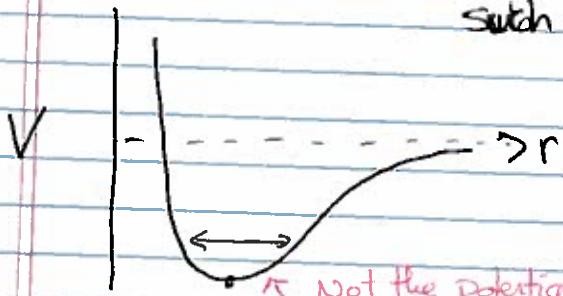
the fundamental frequency

If the molecular potential energy were a perfect harmonic oscillator, the selection rules for v would be

$$\Delta v = \pm 1 \quad \text{IF vibration induces time varying electric dipole}$$

* BUT *

Real molecular potentials have "anharmonicity". This makes it such that larger Δv are electric dipole allowed.



$$\Delta v \geq 1 \text{ allowed}$$

Not the potential well of a perfect harmonic oscillator

$\Delta v \geq 2$ called overtones and usually weaker.

NOTATION

Consider a transition from $V_u \rightarrow V_e$. Depending on how J , the rotational angular momentum, changes, we have different "branches"

| <u>ΔJ</u> | <u>Notation</u> | <u>Means</u> |
|------------------------------|-----------------|----------------------------|
| + | O (J_e) | $J'_{e+2} \rightarrow J_e$ |
| + 2 | P (J_e) | $J'_{e+1} \rightarrow J_e$ |
| + 1 | Q (J_e) | $J'_e \rightarrow J_e$ |
| 0 | R (J_e) | $J'_{e+1} \rightarrow J_e$ |
| - 1 | S (J_e) | $J'_{e+2} \rightarrow J_e$ |
| - 2 | | |
| : | : | |

$V_u \rightarrow V_e$ with corresponding change in J

EXAMPLE 1 : CO

vibration changes μ_e so there are "IR active" ro-vib. transitions

Electric dipole selection rules for rotations $\Delta J = \pm 1$ = we should observe Strong P and R branch transitions

$\Delta J = 0$ forbidden \Rightarrow No Q branch in gas phase

* In solid state (ICES!) there are no rotations so Q branch is a strong band!

Fundamental frequency: γ_0 corresponds to $\lambda_0 \sim 4.6 \mu\text{m}$ in mid-IR Note $E_u(v=1) \sim 3100 \text{ K}$!

Overtone $v=2-0$ close to $\lambda_{2-0} \sim 2.3 \mu\text{m}$ in near-IR

"Hotbands" are any transitions between $V_u, V_e \neq 0$.

Bandheads

Because the R branch and P branch transitions have different sign w.r.t. J , it causes R branch lines to "pile up" and form a bandhead :

Consider $V_u \rightarrow V_e \quad J_u \rightarrow J_e$ ro-vibrational transition

$$\nu_{\text{rovib}} = \nu_0 + B_u J_u (J_u + 1) - B_e J_e (J_e + 1)$$

R branch $J_u = J_e + 1 \quad \Delta J = -1$

$$\nu_R = \nu_0 + B_u (J_u + 1) (J_e + 2) - B_e J_e^2 - B_e J_e$$

$$= \nu_0 + B_u J_e^2 + 3B_u J_e + 2B_u - B_e J_e^2 - B_e J_e$$

$$= \nu_0 + (3B_u - B_e) J_e - (B_e - B_u) J_e^2 + 2B_u$$

↑ Notice that signs are different. Think about why...
At high J_e , the $-x J_e^2$ term will dominate
and ν_R will head back toward ν_0 .

P branch $J_u = J_e - 1 \quad \Delta J = +1$

$$\nu_P = \nu_0 + B_u (J_e - 1) (J_e) - B_e J_e^2 - B_e J_e$$

$$= \nu_0 + B_u J_e^2 - B_u J_e - B_e J_e^2 - B_e J_e$$

$$= \nu_0 - (B_u + B_e) J_e - (B_e - B_u) J_e^2$$

Notice these terms same sign!
So P branch does NOT reverse in frequency at high J_e .

Bandhead is formed by R branch high J lines piling up.